

# Euclid: Early Release Observations – Unveiling the morphology of two Milky Way globular clusters out to their periphery<sup>★</sup>

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## ABSTRACT

As part of the *Euclid* Early Release Observations (ERO) programme, we analyse deep, wide-field imaging from the VIS and NISP instruments of two Milky Way globular clusters (GCs), namely NGC 6254 (M10) and NGC 6397, to look for observational evidence of their dynamical interaction with the Milky Way. We search for such an interaction in the form of structural and morphological features in the clusters' outermost regions, which are suggestive of the development of tidal tails on scales larger than those sampled by the ERO programme. Our multi-band photometric analysis results in deep and well-behaved colour–magnitude diagrams that, in turn, enable an accurate membership selection. The surface brightness profiles built from these samples of member stars are the deepest ever obtained for these two Milky Way GCs, reaching down to  $\sim 30.0$  mag arcsec<sup>-2</sup>, which is about 1.5 mag arcsec<sup>-2</sup> below the current limit. The investigation of the two-dimensional density map of NGC 6254 reveals an elongated morphology of the cluster peripheries in the direction and with the amplitude predicted by *N*-body simulations of the cluster's dynamical evolution, at high statistical significance. We interpret this as strong evidence for the first detection of tidally induced morphological distortion around this cluster. The density map of NGC 6397 reveals a slightly elliptical morphology, in agreement with previous studies, which requires further investigation on larger scales to be properly interpreted. This ERO project thus demonstrates the power of *Euclid* in studying the outer regions of GCs at an unprecedented level of detail, thanks to the combination of large field of view, high spatial resolution, and depth enabled by the telescope. Our results highlight the future *Euclid* survey as the ideal data set to investigate GC tidal tails and stellar streams.

**Key words.** Galaxy: evolution – globular clusters: general – Galaxy: structure – techniques: photometric – Stars: imaging

## 1. Introduction

Globular clusters (GCs) have long played an important role as cosmic laboratories for a wealth of astronomical open questions. As the closest example in nature of simple stellar populations,

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they are exceptionally well suited to investigate stellar evolution (Salaris et al. 1993; Cassisi et al. 1999). Due to their short dynamical timescales, GCs are ideal for studying dynamical processes and their interplay with stellar evolution itself (Hut et al. 1992; Meylan & Heggie 1997). Moreover, being some of the oldest objects in the Universe, they are witnesses and powerful tracers of the early formation and evolution of galaxies like the Milky Way (Searle & Zinn 1978; Forbes & Bridges 2010; Massari et al. 2019; Kruijssen et al. 2020; Massari et al. 2023).

A major challenge when studying GCs arises from their richness in stars. Especially in the inner regions, their high stellar density causes crowding to be severe, so much so that these objects have been observed mostly with high-resolution, limited field-of-view (FoV) cameras, which have sufficient depth to sample their stellar component down to the faintest, lowest-mass stars. A clear example of this is provided by the success that the *Hubble* Space Telescope (HST) has demonstrated in the investigation of GCs (e.g., Sarajedini et al. 2007; Piotto et al. 2015), and in particular of their central regions. However, the information provided by low-mass stars in GC outskirts is crucial for understanding how these systems evolve both internally and in the presence of the Milky Way.

For example, it is well known that any stellar system embedded in a gravitational potential is subject to tidal forces (Binney & Tremaine 1987). This effect is more pronounced where the tidal forces are comparable to, or start to dominate over the system’s self-gravity, that is in the outer peripheries. Stars in the outskirts of GCs are thus more prone to have their orbits distorted by tidal forces, contributing to the development of a cluster tidal tail. Since the internal dynamical evolution of GCs drives faint, low-mass stars to preferentially populate the cluster outer regions, GC tidal tails appear as morphological features having very low surface brightness (Grillmair et al. 1995; Leon et al. 2000).

The detection of tidal tails around GCs has proven challenging so far. On the one hand, dense GCs experience low evaporation rates during their evolution, so that stars in their tidal tails typically constitute less than 0.1% of the surrounding stellar field (Dinescu et al. 1999; Balbinot & Gieles 2018; Sollima 2020). On the other hand, tidal tails emerge over scales that can be several to tens of degrees wide (Odenkirchen et al. 2001; Grillmair & Dionatos 2006; Erkal et al. 2017), making them difficult to sample with existing high-resolution imagers. The advent of photometric surveys, such as SDSS (York et al. 2000), Gaia (Gaia Collaboration et al. 2016), and DES (Dark Energy Survey Collaboration et al. 2016), has at least partially solved the latter of these issues, and led to the detection of tidal tails around several GCs (e.g., Odenkirchen et al. 2001; Belokurov et al. 2006; Niederste-Ostholt et al. 2010; Sollima et al. 2011; Balbinot et al. 2011; Myeong et al. 2017; Navarrete et al. 2017; Ibata et al. 2019; Grillmair 2019; Kaderali et al. 2019; Carballo-Bello 2019; Bianchini et al. 2019; Shipp et al. 2020; Piatti et al. 2021). Yet, the number of systems displaying tidal tails remains small (Kuzma et al. 2018; Sollima 2020). The absence of tails might be explained by the presence of dark matter halos surrounding GCs, which would prevent stars from escaping (Moore 1996), and would favour the development of diffuse stellar envelopes (Peñarrubia et al. 2017; Kuzma et al. 2018). Alternatively, the reason why tails remain such an elusive feature could be an observational bias, related to the fact that existing imagers are not able to observe very faint stars on large scales on the sky (Balbinot & Gieles 2018).

Thanks to its wide-field imaging capabilities, the *Euclid* mission (Laurijs et al. 2011; Euclid Collaboration: Mellier et al.

2024) will provide a definitive answer to this fundamental open question concerning the nature of GCs. The Euclid Wide Survey will observe an area of approximately 14 000 deg<sup>2</sup> with exceptionally deep AB limiting magnitudes ( $5\sigma$  point-like sources) of 26.2 in the optical VIS band ( $I_E$ ) and 24.5 in the near-infrared NISP bands ( $Y_E$ ,  $J_E$ ,  $H_E$ , see Euclid Collaboration: Scaramella et al. 2022). This means that all of the GCs covered by the survey footprint will be observed with an unprecedented combination of depth, spatial coverage, and resolution, in four different photometric bands.

In this paper, we showcase the power of *Euclid* for the investigation of Milky Way GC peripheries by analysing Early Release Observations (ERO, Euclid Early Release Observations 2024) of two GCs, namely NGC 6254 (M10) and NGC 6397. NGC 6254 has already been the subject of some wide-field photometric studies. In the first of these, Leon et al. (2000) analysed photometric plates taken with the ESO Schmidt telescope, covering a  $5.5 \times 5.5$  deg<sup>2</sup> area around the cluster and having a limiting magnitude of  $R \simeq 19$ , thus barely reaching the cluster main-sequence turn-off. The authors found the presence of tidal tails extending along the north–south direction, but they highlighted a possible bias in the detection induced by the strong gradient in the dust extinction across the field. In a later study, Dalessandro et al. (2013a) investigated the density profile of the cluster using deeper ( $V \simeq 20$ ) photometry obtained from observations taken with the Wide-Field Imager mounted at the 2.2-m ESO-MPG telescope over a  $33 \times 34$  arcmin<sup>2</sup> FoV. In this case, the authors found a density profile that was well fit by a simple King model (King 1962), thus possibly indicating a lack of asymmetric morphological features like tidal tails. According to the predictive algorithm by Balbinot & Gieles (2018), which takes into account the mass loss experienced by a GC and its orbital phase, NGC 6254 has a high chance of having tidal tails as it is very close to its apocentre (orbital phase  $\Phi = 0.91$ ) and it should have already lost 61% of its initial mass.

The detection of tidal tails around NGC 6397 is debated, too. The first to look for extra-tidal features were Leon et al. (2000), who found overdensities that were classified as unreliable, due to the uncertainty in the distribution of dust around the cluster. More recently, as Gaia-related results, Ibata et al. (2021, 2023) found the presence of a possible tail extending more than 18 degrees on the sky, but Boldrini & Vitral (2021) challenged these findings by not detecting any tail. According to the metric by Balbinot & Gieles (2018), NGC 6397 is a strong candidate for having developed tidal tails, as it has lost 72% of its initial mass and its orbital phase is  $\Phi = 0.95$ .

As part of the *Euclid* ERO programme, the available imaging covers an almost square region of about  $0.8 \times 0.8$  deg<sup>2</sup> centred on the GCs. At a distance of 5.07 kpc (Baumgardt & Vasiliev 2021), the tidal radius of NGC 6254 is about  $r_t \sim 0.3$  deg (Dalessandro et al. 2013a), and thus will be entirely covered by the observations. The spatial coverage is more limited for NGC 6397, whose tidal radius is larger, being  $r_t \sim 0.6$  deg (Moreno et al. 2014), at a distance of 2.48 kpc (Baumgardt & Vasiliev 2021). This means that it is likely these ERO data will not enable the detection of tidal tails, which start to develop well outside the tidal radius, at the location of the Jacobi radius (0.6 deg for NGC 6254 and 1.15 deg for NGC 6397, see Webb et al. 2013). Nonetheless, they will give us an unprecedented chance to detect the morphological distortions in the clusters’ outer regions, which may ultimately hint at the presence of tidal tails on larger scales.

This paper is structured as follows. In Sect. 2, the observations of the two GCs are described. In Sect. 3, we present the results of the photometric analysis by showing the first *Euclid*

colour–magnitude diagrams (CMDs, the magnitudes are in the AB photometric system) of Milky Way GCs. The morphology of the GCs’ outer regions is discussed in Sect. 4, and it is further interpreted in light of  $N$ -body simulations in Sect. 5. Finally, the conclusions of our work are summarised in Sect. 6.

## 2. Early Release Observations of NGC 6254 and NGC 6397

The observations of NGC 6254 and NGC 6397 have been performed on the 9th and the 22nd of September 2023, respectively, as part of the *Euclid* ERO Programme. This programme (which is described in detail in Cuillandre et al. 2024a, C24 hereafter) further includes observations of the  $\sigma$  Orionis cluster (Martin et al. 2024), of nearby galaxies (Hunt et al. 2024), of the Fornax galaxy cluster (Saifollahi et al. 2024), of the Perseus cluster of galaxies (Cuillandre et al. 2024b; Kluge et al. 2024; Marleau et al. 2024) and of a giant gravitational lens (Atek et al. 2024; Weaver et al. 2024). Each observation analysed here consists of a single reference observation sequence (ROS, see *Euclid Collaboration: Scaramella et al. 2022*) centred on the cluster. One ROS includes four different dithered pointings, each resulting in one 560 seconds-long exposure in the  $I_E$  band (*Euclid Collaboration: Cropper et al. 2024*), and in 87.2 seconds-long exposures in the  $Y_E$ ,  $J_E$  and  $H_E$  bands<sup>1</sup>, plus one spectral exposure that is neglected for this study. All of the exposures have been reduced by means of the data reduction pipeline described by C24. In the following, we briefly describe the main outputs of the pipeline.

### 2.1. VIS and NISP imaging

Figure 1 shows the distribution on the plane of the sky of the four deep photometric exposures of NGC 6254 ROS in the  $I_E$  and NISP bands. The VIS exposures have been stacked together following the prescriptions in C24, achieving an internal  $1\sigma$  astrometric precision of  $\sim 4$  mas. The same stacking procedures of C24 have been applied on the NISP (*Euclid Collaboration: Schirmer et al. 2022; Euclid Collaboration: Jahnke et al. 2024*) exposures as well. In this case the astrometric referencing onto Gaia DR3 positions achieved a  $1\sigma$  precision of 15 mas, which is still less than one tenth of the pixel scale ( $0.3$  arcsec pixel<sup>-1</sup>, see *Euclid Collaboration: Schirmer et al. 2022*).

The measured full width at half maximum (FWHM) in the stacked images is typically 0.16 arcsec in the VIS band and  $\sim 0.4$  arcsec in the NISP ones, and is stable across the FoV within a few percent.

The dither pattern allowed for a continuous coverage of the field around the clusters. The large FoV coverage and high spatial resolution are exceptionally represented by, respectively, the left- and the right-hand panels of Fig. 2, which show an RGB image of the field around NGC 6397. The FoV coverage of NGC 6254 resulted in a slightly imperfect gap filling due to the dither pattern. This does not impact our scientific objectives.

## 3. Photometric analysis

The stacked images described above have been analysed using the Astromatic software suite, including SExtractor and PSFEx (Bertin 2011). Briefly, the PSFEx code models the point

<sup>1</sup> The quoted numbers refer to the effective exposure times, whereas the total duration indicated in *Euclid Collaboration: Scaramella et al. (2022)* includes overheads.

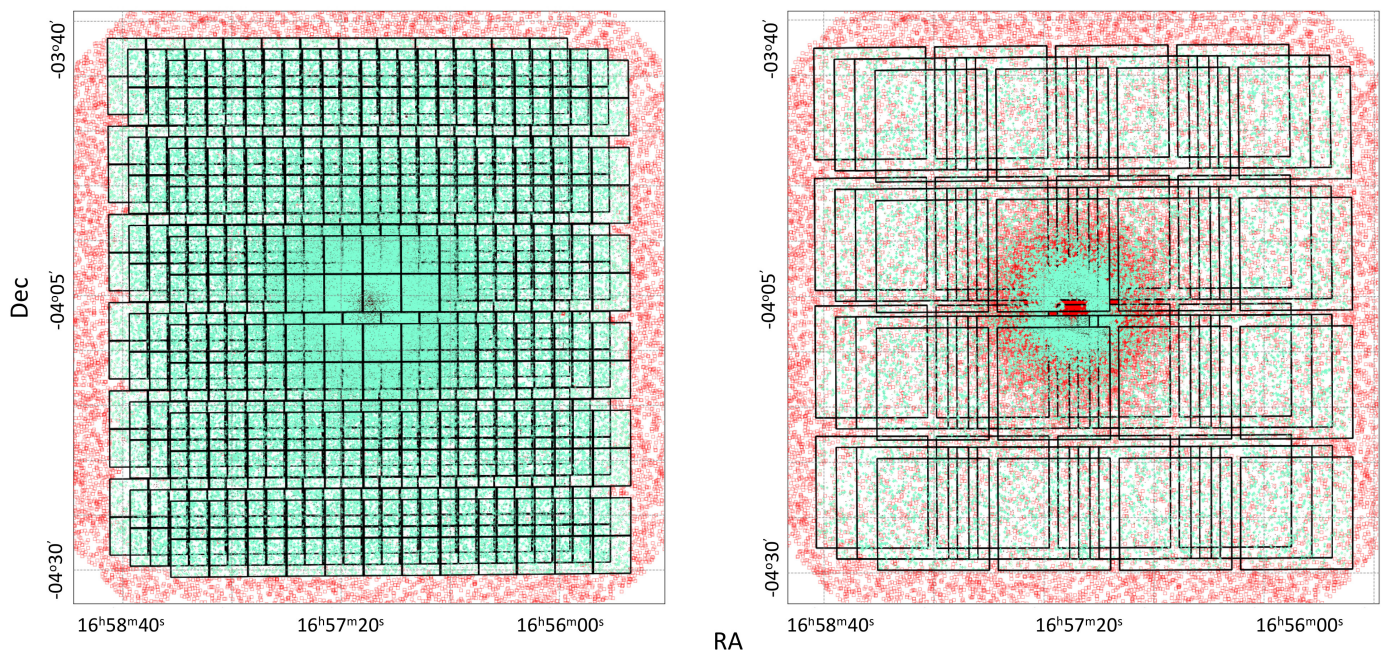
spread function (PSF) of an instrument as a combination of basis functions and works directly on the image. In particular, PSFEx selects point-like sources detected by SExtractor, and uses them to compute the PSF model. For both the GCs investigated in this work, we selected stars for the PSF model that are bright, non-saturated, and relatively isolated in each band; specifically, we only selected point-like sources (with a maximum ellipticity of 0.3 and a FWHM within 30% of the nominal one) having at least 5 pixels above a detection threshold of  $10\sigma$  above the background. A further filtering through a Moffat function has been imposed to reject spurious detections. The PSF model has been determined over a  $6 \times 6$  equally spaced grid covering the whole FoV, allowing for a linear spatial variation. In this way, for both the GCs, the number of sources used to model the PSF exceeded a few thousand. The final PSF models have an average ellipticity of 0.02 and a FWHM spatial variation of only 3%.

These models have then been used by SExtractor to perform PSF photometry on a much larger sample of sources, defined by relaxing the source-detection threshold to having even a single pixel  $1.5\sigma$  above the background. In this way, we PSF-fit more than 400 000 sources in the VIS image of NGC 6254, and more than one million sources in the VIS image of NGC 6397. The number of detections in the NISP images is obviously lower because of reduced depth and resolution, about 190 000 for NGC 6254 and about 350 000 for NGC 6397. By combining the VIS and NISP catalogues, most of the artefacts and spurious detections (e.g., PSF spikes around very bright stars) are rejected, and the final multi-band catalogues for NGC 6254 and NGC 6397 contain positions, magnitudes, related uncertainties, and shape/quality parameters for 133 090 and 296 646 sources, respectively. Aperture corrections of a few percent (consistent with other works on ERO data, see e.g., Saifollahi et al. 2024), estimated using bright and isolated stars, are applied to the measured magnitudes. Because of the very high crowding, only a modest number of stars have been fit in the core of the clusters, where incompleteness is therefore very high. We remark, though, that throughout this paper we do not focus on the clusters’ central regions, and thus our conclusions are not affected by this. A quantitative assessment of the local completeness achieved by the ERO data of NGC 6254 in the VIS band is provided in the next section.

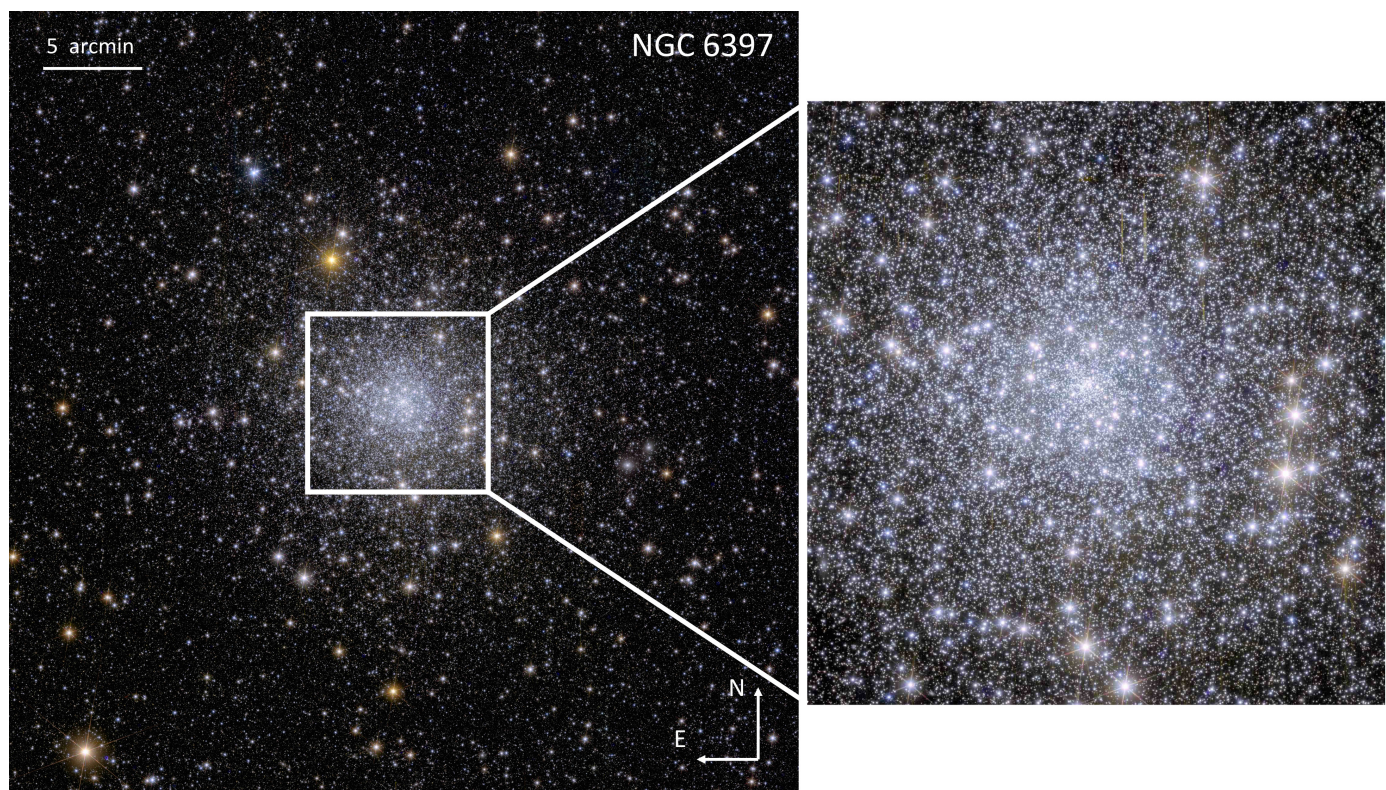
### 3.1. Completeness: local estimate for NGC 6254

Directly estimating the completeness of the photometry would require artificial-star tests that are computationally very expensive for images this large in size, and that are not yet available in the adopted data-reduction pipeline. Nonetheless, we can take advantage of the particularly well suited, publicly available photometric catalogue by Dalessandro et al. (2011). These authors analysed HST Wide Field Planetary Camera 2 (WFPC2) observations in the F606W and F814W bands of a stellar field located between 1 and  $2 r_h$  of the cluster centre ( $r_h = 2.02$  arcmin according to Baumgardt & Vasiliev 2021). By means of artificial-star tests performed in Beccari et al. (2010), the authors computed the completeness of their WFPC2 photometric catalogue, which we can therefore use for a relative comparison.

We cross-matched our VIS catalogue with the public photometry of Dalessandro et al. (2011) by means of a linear transformation, requiring matching sources to (a) be located within a distance of 0.5 arcsec and (b) have a difference between the *Euclid*  $I_E$  and the WFPC2  $m_{F814W}$  magnitudes smaller than a conservative value of 2 mag. The residuals from the resulting linear transformation solution are only 0.03 arcsec in both Right As-



**Fig. 1.** Sky coverage and dither pattern of the four VIS (left-hand panel) and NISP (right-hand panel) exposures of the ERO ROS for NGC 6254. Red symbols indicate *Gaia* reference stars. Green symbols represent *Euclid* detections.



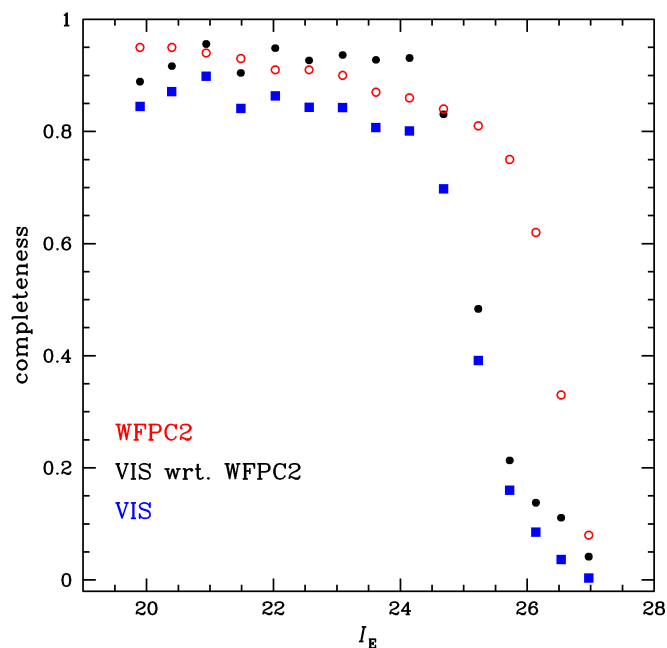
**Fig. 2.** RGB image of NGC 6397. The left-hand panel shows the entire FoV covered by the observations. The right-hand panel shows a  $10 \times 10$  arcmin<sup>2</sup> zoomed-in detail in the GC core. Credits: ESA/*Euclid*/*Euclid* Consortium/NASA, image processing by J.-C. Cuillandre (CEA Paris-Saclay), G. Anselmi, CC BY-SA 3.0 IGO.

cension (RA) and Declination (Dec), thus ensuring an optimally performed cross-match.

Figure 3 shows the results of the comparison. We remark that, in order to perform a meaningful analysis, we decided to exclude regions around very bright, saturated stars, which are populated by artefacts and PSF structures in both the *Euclid* and

HST images. The red open symbols in Fig. 3 show the absolute completeness as a function of magnitude for the WFPC2 sample, as estimated by Dalessandro et al. (2011). The  $m_{F814W}$  magnitudes were converted to  $I_E$  magnitudes by means of a linear transformation depending on  $m_{F555W} - m_{F814W}$  colour, determined by using the stars in common between the two cata-

logues. On the other hand, the black filled symbols indicate the *relative* completeness of the VIS catalogue with respect to the WFPC2 one. As shown in the Figure, the trend remains near-constant around a value of 94% down to  $I_E = 24$ , and then drops to lower percentages. The product of these two curves, shown as



**Fig. 3.** Completeness of the *Euclid* ERO in the  $I_E$  band, at a distance of  $1.5 r_h$  from the core of NGC 6254. Red empty circles indicate the completeness of the reference HST/WFPC2 catalogue. Black circles represent the relative completeness of *Euclid* with respect to HST. Finally, the filled blue squares show the absolute completeness of the VIS data.

blue filled squares in Fig. 3, provides an estimate for the *absolute* completeness of the VIS catalogue, at an average distance from the GC centre of  $1.5 r_h$ . This is significantly closer to the most crowded regions of the GC than the peripheries we are interested in. Moreover, this estimate does not take into account possible VIS detections that are missing in the WFPC2 catalogue. For these reasons, it is safe to consider this estimate of the catalogue completeness as a *lower limit*. Completeness is certainly better at larger distances in our data. As can be seen, the completeness remains about constant at  $\sim 82\%$  down to  $I_E = 24$ , dropping to 50% completeness by  $I_E \simeq 25.0$ .

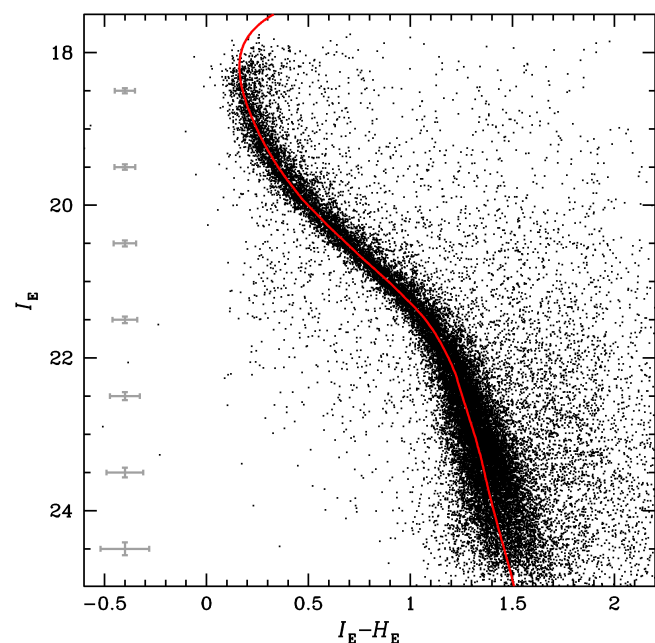
The analysis of the VIS catalogue completeness has been performed to understand the ERO data better and to quantify the performance achieved, but it is not used for the study presented in this paper. These findings are of fundamental importance for future analysis of these ERO data, though, as a wealth of science cases, such as the investigation of the cluster mass functions, or of their binary populations, requires the knowledge of the photometric completeness and its correction.

### 3.2. Colour–magnitude diagrams

A fundamental quality check on the photometric analysis of GC fields comes from the cluster colour–magnitude diagram (CMD). In the following, we describe the most important features we found in the CMDs of NGC 6254 and NGC 6397, which are the first *Euclid* CMDs of Milky Way stellar systems ever presented.

#### 3.2.1. NGC 6254

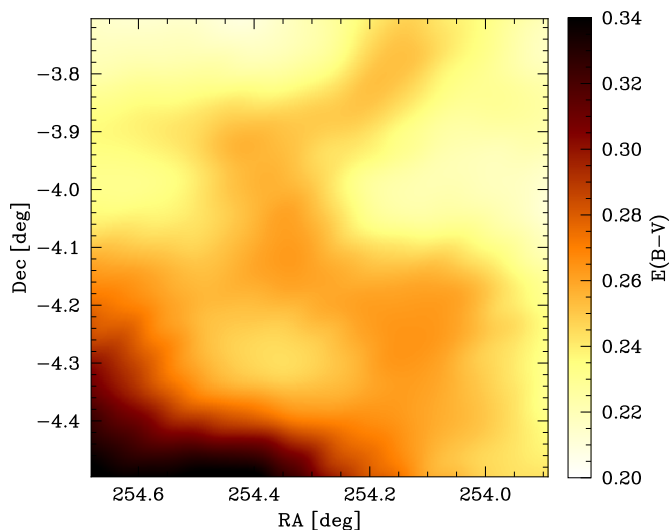
Figure 4 shows the  $(I_E - H_E, I_E)$  CMD for an annulus between 5 and 10 arcmin from the cluster centre, as determined by Dalesandro et al. (2013a). This region has been selected as it is close enough to the cluster for the CMD to be dominated by GC stars, but far enough (well beyond  $2 r_h$ ) to avoid severe incompleteness and crowding effects. The photometry has been corrected for differential reddening by using the colour excess  $E(B - V)$  map from Schlegel et al. (1998) evaluated using *mwdust* (Bovy et al. 2016) with the corrections from Schlafly & Finkbeiner (2011) and the extinction law by Cardelli et al. (1989). We correct all magnitudes to a fiducial value of  $E(B - V) = 0.28$ , the mean absolute colour excess of the cluster (Harris 1996, see Fig. 5). To improve upon the correction provided by the quoted map, we further apply the differential reddening correction procedure described by Milone et al. (2012). By design, such an additional step effectively corrects also for local residuals in the photometric calibration, which can manifest as slight systematic differences in colour. Indeed, we find additional corrections of the order of a few hundredth mag, which we have applied to the observed magnitudes in Fig. 4. Overall, the CMD shows a well-defined main sequence (MS), which extends from the turn-off (MSTO) point at  $I_E \approx 18.5$  mag (where the deep exposures start to saturate) to a couple of magnitudes below the MS knee (see e.g. Bono et al. 2010; Massari et al. 2016; Saracino et al. 2018), located at about  $I_E = 21.5$  mag.



**Fig. 4.** Differential reddening corrected,  $(I_E - H_E, H_E)$  CMD for an annulus between 300 arcsec and 600 arcsec from the centre of NGC 6254. Error bars are marked in grey. A 13 Gyr old,  $[\text{Fe}/\text{H}] = -1.5$  dex and  $\alpha$ -enhanced theoretical isochrone from the BaSTI archive is shown in red for the sake of comparison.

A direct comparison with a theoretical isochrone is also shown. The isochrone has been taken from the BaSTI archive (Pietrinferni et al. 2021) and describes a 13 Gyr old population having  $[\text{Fe}/\text{H}] = -1.5$  dex and  $\alpha$ -enhanced ( $[\alpha/\text{Fe}] = 0.4$ ) chemical composition (Haynes et al. 2008). To take into account possible uncertainties on the absolute photometric calibration, as well as on the adopted reddening estimates (Harris 1996), the

isochrone has been shifted in colour by a small arbitrary amount ( $\delta[I_E - H_E] = 0.02$ ) to provide a satisfactory fit. Once the shift is applied, the model describes the observed MS very well and shows that the observed CMD samples masses down to  $0.16 M_\odot$  (at  $I_E \approx 24.9$ ). This is the deepest CMD of a GC over such a large scale.



**Fig. 5.** Extinction map around NGC 6254 as obtained from the Schlafly & Finkbeiner (2011) corrections to the extinction map by Schlegel et al. (1998).

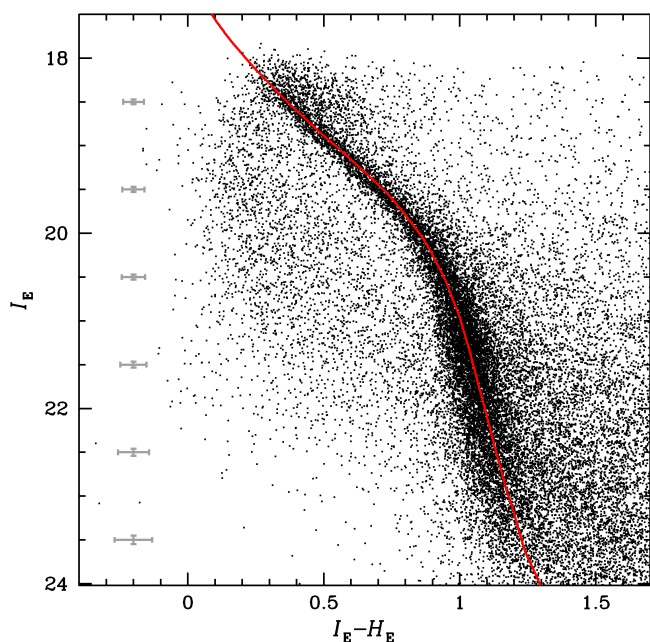
Another interesting feature that can be appreciated from this CMD is the possible widening of the faint MS below the MS knee. From a statistical point of view, the observed colour dispersion of upper-MS stars in the bright, non-saturated magnitude range  $18.75 < I_E < 19.25$  is  $\sigma_{\text{obs}}^{\text{uMS}} = 0.04$ . The typical nominal colour error at those magnitudes, as estimated by SExtractor, is of only 0.01 mag. Given that NGC 6254 is not a peculiar type-2 GC (Milone et al. 2019), and therefore should not present photometric features related to large iron or helium dispersion, which could significantly widen the observed sequences, such a nominal error seems to underestimate the actual one. We thus take the observed standard deviation  $\sigma_{\text{obs}}^{\text{uMS}}$  of GC members at these bright magnitudes as a more robust (even though conservative) estimate of the photometric colour uncertainty. Possible terms that could inflate this value, making it an overestimate of the actual error, are the presence of photometric binaries, small helium variations among cluster stars, and residuals in the differential reddening correction. The observed binary fraction of NGC 6254 at the distance sampled by the CMD shown in Fig. 4 is of only 1.5% (Dalessandro et al. 2011), and therefore should not contribute significantly to  $\sigma_{\text{obs}}^{\text{uMS}}$ . According to Milone et al. (2018), NGC 6254 might be characterised by a small helium spread  $\Delta Y < 0.029$ , that could contribute to a photometric spread of 0.01–0.02 mag.

This colour dispersion should also increase for fainter stars due to the decrease in the signal-to-noise ratio (SNR), according to  $\sigma_{\text{obs}}^{\text{uMS}} \approx \log_{10}(1 + 1/\text{SNR})$ . Assuming a systematic colour uncertainty floor of  $\sigma_{\text{obs}}^{\text{uMS}} = 0.04$  at  $I_E = 19$ , we thus expect  $\sigma_{\text{obs}}^{\text{uMS}} = 0.09$  below the MS knee, in the range  $23 < I_E < 23.5$  (see the grey error bars in Fig. 4, where the error in magnitude has been computed by assuming an equal contribution to  $\sigma_{\text{obs}}$  from the two bands). However, the observed colour dispersion at the same magnitudes is somewhat larger, reaching  $\sigma_{\text{obs}} = 0.12$ . Given our conservative colour error estimate, we tentatively in-

terpret this widening as an intrinsic feature, and cautiously attribute it to the GC’s multiple stellar populations (see e.g., Gratton et al. 2004; Milone et al. 2017), which manifest in near-IR CMDs as a widening/splitting of the MS below the MS-knee due to the opacity effect of collision-induced absorption by water on the surface of M-dwarfs (Milone et al. 2012, 2019). The presence of multiple-populations in NGC 6254 is widely known, both from photometric (Monelli et al. 2013; Milone et al. 2017) and spectroscopic (Carretta et al. 2009) studies. Yet, future more detailed analysis of this CMD will test the significance of this feature on more solid statistical ground, thus assessing the effectiveness of *Euclid* in investigating this peculiar property of GCs.

### 3.2.2. NGC 6397

Figure 6 shows the  $(I_E - H_E, I_E)$  CMD for an annulus at distances  $r \in [300, 600]$  arcsec from the centre of NGC 6397. The CMD is corrected for differential reddening and residuals on the photometric calibration as described for NGC 6254.



**Fig. 6.** Differential reddening corrected,  $(I_E - H_E, I_E)$  CMD for an annulus between 300 arcsec and 600 arcsec from the centre of NGC 6397. Error bars are marked in grey. A 13 Gyr old,  $[\text{Fe}/\text{H}] = -2.0$  dex, and  $\alpha$ -enhanced theoretical isochrone from the BaSTI archive is shown in red for the sake of comparison. The effect of saturation becomes evident at  $I_E \lesssim 18.5$ .

Also for NGC 6397, the cluster MS stands out very clearly as a well-defined and populated sequence, even though the contamination by field stars and background galaxies appears stronger than in the case of NGC 6254. The higher density of sources is likely the cause for the shallower depth reached by the CMD. The faintest end of the cluster MS seems to reach a magnitude of  $I_E \sim 24.0$ , about 1 mag brighter than in the case of NGC 6254. Such a shallower cutoff is imposed by the limiting magnitudes in the  $H_E$  band detections, that suffer from the low spatial resolution in such a crowded field, and from persistence effects.

Because of the closer distance of NGC 6397 to us, saturation prevents photometric measurement of the MSTO, but we are able to sample even lower stellar masses than in the

case of NGC 6254. According to the BaSTI isochrone overplotted in Fig. 6, which represents a 13 Gyr old population with  $[\text{Fe}/\text{H}] = -2.0$  dex and  $\alpha$ -enhanced chemical composition, the MS sampled lower limit corresponds to  $\sim 0.12 M_{\odot}$ .

The colour dispersion around the non-saturated upper MS (at  $I_E \lesssim 18.5$  mag) gives once again a reasonable upper limit to the achieved photometric error of  $\sigma_{\text{obs}}^{\text{uMS}} = 0.04$  mag, very similar to that found for NGC 6254. Following the same argument as in that cluster, the colour error below the MS knee, at  $22 < I_E < 22.5$  should increase to 0.06 for NGC 6397 (see the error bars in Fig. 6). Again, we detect a marginal widening of the faint MS ( $\sigma_{\text{obs}} = 0.08$ ), possibly ascribable to the already known phenomenon of multiple populations in this cluster (see e.g., di Criscienzo et al. 2010; Milone et al. 2018; Carretta et al. 2009).

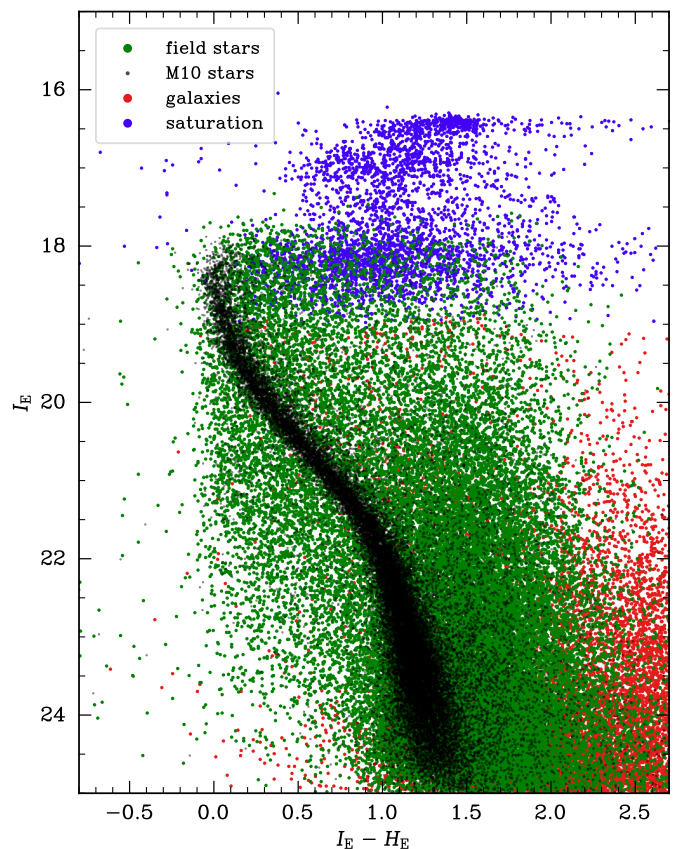
### 3.3. Star-galaxy and cluster-field separation

In order to study the morphology of the outer regions of the two GCs, it is fundamental to separate likely member GC stars from contaminating background galaxies and Milky Way disc stars. In the absence of kinematic information for all the individual sources, the best way to isolate cluster stars is to take advantage of the multi-band, high-resolution capabilities of *Euclid* and to combine colour information with photometric parameters that are sensitive to the point-like or extended shape of each source. Figure 7 shows the  $(I_E - H_E, I_E)$  CMD of all sources falling in the ERO FoV around NGC 6254. When compared to the CMD shown in Fig. 4, the much higher complexity of the sampled stellar population becomes evident. In order to unravel such a complexity, we make use of two of the fit parameters provided by *SExtractor*, called SPREAD\_MODEL and CLASS\_STAR. The former parameter is designed<sup>2</sup> such that it equals zero for stars, and deviates from zero for more extended (or more compact) sources. The latter parameter, instead, equals 1 for sources that are well fit by the adopted PSF model, and is  $< 1$  for sources with a less satisfactory fit. Figure 8 shows the behaviour of these parameters as a function of  $I_E$  magnitude. It is immediately evident that stars (yellow symbols with CLASS\_STAR  $\sim 1$ ) can be rather well distinguished from other sources by imposing a selection on  $|\text{SPREAD\_MODEL}| < 0.01$ . What is left out of this selection corresponds either to background galaxies (if CLASS\_STAR is  $\ll 1$ , see the red symbols in Fig. 7) or to saturated stars (if CLASS\_STAR  $\sim 1$ , see blue symbols in Fig. 7).

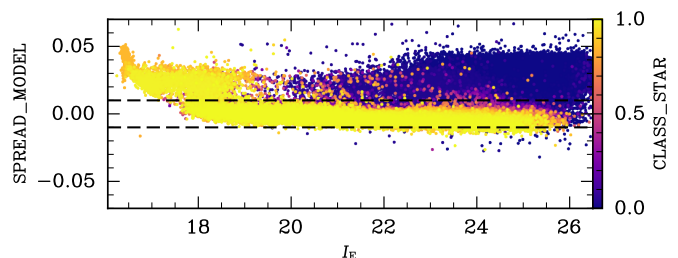
Among what is selected as point-like sources, the stars located closer to the GC centre, and shown as black symbols in Fig. 7, describe the location of cluster members, whereas the remaining stars, shown as green symbols, represent the numerous Milky Way field stars, composed of a mix of distant halo MS stars and nearby M-dwarfs. In this case, the colour information plays the crucial role, and in particular we find that the longest colour baseline ( $I_E - H_E$ ) is the combination of filters that best separates field stars, galaxies, and the GC stellar populations.

In summary, the CMD analysis performed in this section on NGC 6254 provides the ingredients to select likely member GC stars in the best way possible, in the absence of kinematic information. The use of the same SPREAD\_MODEL and CLASS\_STAR parameters, together with a colour–magnitude selection has been effective in the case of NGC 6397 as well, and is adopted hereafter, unless specifically indicated.

<sup>2</sup> see <https://sextractor.readthedocs.io/en/latest/Model.html>



**Fig. 7.** CMD showing likely members of NGC 6254 (black), field stars (green), and background galaxies (red).



**Fig. 8.** SPREAD\_MODEL as a function of  $I_E$  colour-coded by CLASS\_STAR. The dashed lines indicate a selection of  $|\text{SPREAD\_MODEL}| < 0.01$ .

## 4. Results

### 4.1. Surface-brightness profiles

In order to build a GC surface-brightness profile that is sensitive even to very low stellar surface density, we select a sample of likely cluster member stars by using the photometric parameters described before and by targeting completeness over purity. In this sense, we applied only conservative cuts of  $\text{CLASS\_STAR}_{I_E} > 0.03$ ,  $|\text{SPREAD\_MODEL}|_{I_E} < 0.01$ ,  $|\text{SPREAD\_MODEL}|_{H_E} < 0.01$ , and  $\text{FLUX\_RADIUS}_{I_E} < 1.5$ ,<sup>3</sup> to exclude obvious artefacts and extended sources. Then, we select our sample of GC stars along a region in the reddening-corrected CMD centred around the theoretical isochrone shown in Fig. 4, down to  $I_E = 24$ , and in a band

<sup>3</sup> FLUX\_RADIUS is a *SExtractor* parameter that describes the half-light radius of a source light profile.

$\sim 0.3$  mag wide in  $I_E - H_E$  colour. As described by Fig. 7, this selection should include the vast majority of GC members.

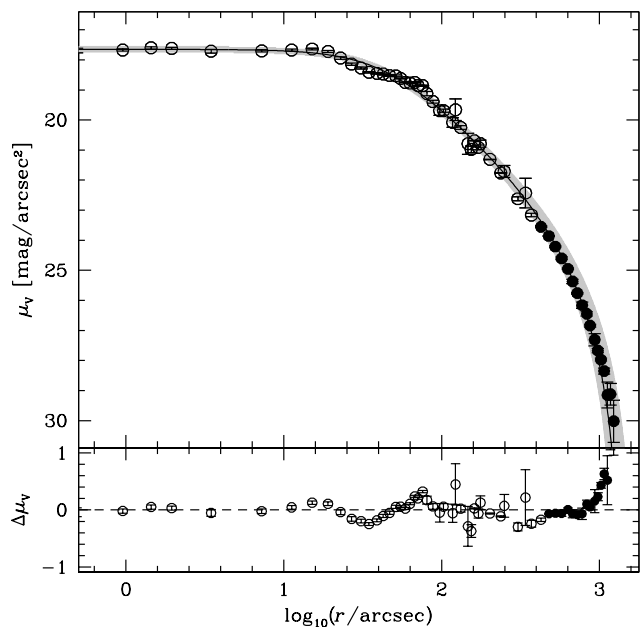
By using these samples of selected stars, we obtain the surface-density profiles from resolved star counts for both clusters. For NGC 6254, we sub-divided the total FoV in 25 concentric annuli centred on the cluster’s centre of gravity from [Dalessandro et al. \(2013a\)](#), reaching a maximum cluster-centric distance  $r = 1400''$ , while we split the area sampled for NGC 6397 into 28 annuli centred on the centre reported by [Goldsbury et al. \(2010\)](#), thus reaching a maximum radial coverage of  $r = 1800''$ . Each annulus was then split in an adequate number of angular sub-sectors (two or four depending on the number of stars). In each sub-sector, the density has been estimated as the ratio between the selected number of stars counted and the covered area. The density assigned to a given annulus is the average of the densities of each sub-sector of that annulus. The error assigned to each density measure is defined as the dispersion around the mean of the individual sub-sector densities. We estimated the background density by using the most external data points, which attained almost constant density values. For NGC 6254 we used the six most external density values located at  $r > 1100''$ , while for NGC 6397 we could only use two data-points at  $r > 1600''$ . It is worthwhile noticing that, especially in the latter case, a few GC members are still present at these extreme radii, hence the background density might be slightly over-estimated. We then subtracted the mean background values from all the other annuli. Since the *Euclid* data are strongly incomplete in the inner cluster regions, we combined our resolved star count profiles with the integrated light surface brightness profiles used by [Baumgardt & Hilker \(2018\)](#). To this aim, we first reported our density profiles to the same scale and photometric reference systems of the literature profiles, by using values in common in the radial range  $250'' < r < 400''$  for NGC 6254 and  $400'' < r < 650''$  for NGC 6397. The resulting combined surface density profiles are shown in Fig. 9 and Fig. 10 for NGC 6254 and NGC 6397, respectively.

The first key result coming from this analysis is that, while sampling almost the same radial range as previous studies, we reach surface brightness values  $\mu_V \sim 30.0$  mag arcsec $^{-2}$  for NGC 6254 and  $\mu_V \sim 28$  mag arcsec $^{-2}$  for NGC 6397, which is more than one magnitude fainter than previously obtained (e.g., [Baumgardt & Hilker 2018](#)). This performance already demonstrates the advance that *Euclid* will guarantee on this kind of study in the future.

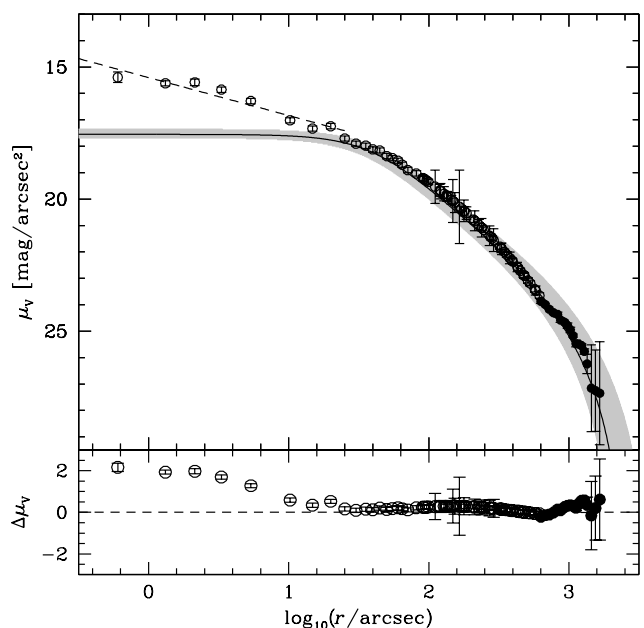
We performed a fit of the azimuthally averaged density profiles by using single-mass King models ([King 1962](#)) for both clusters, and show the results of the fit in Fig. 9 and Fig. 10. The general approach follows the one described in [Dalessandro et al. \(2013b\)](#). For NGC 6254 the best fit is obtained with a core radius  $r_c = 44.2^{+1.7}_{-1.3}$  arcsec and a concentration parameter  $c = 1.47^{+0.07}_{-0.05}$ , which yields a value for the tidal radius  $r_t = 1375.7^{+216.2}_{-247.3}$  arcsec. These results are in good agreement within the errors with previous estimates in the literature (e.g., [Harris 1996](#); [Dalessandro et al. 2013a](#)). We note that  $r_t$  is  $\sim 20\%$  larger (but still compatible within the errors) than in previous derivations. Such a large value can be naturally explained by our use of faint, low-mass MS stars, which tend to preferentially populate the cluster outer regions, while results in the literature are driven by the distribution of heavier red giant branch stars, which are generally more concentrated towards the cluster centre.

For NGC 6397, we excluded from the fit the eight innermost points ( $r < 12''$ ) of the observed profile, since they appear to deviate from a flat core behaviour, as expected for a post-core collapse cluster. In this way, the best-fit model is obtained for

a core radius  $r_c = 42.2^{+2.5}_{-5.3}$  arcsec, concentration  $c = 1.77^{+0.16}_{-0.12}$  and  $r_t = 2580.1^{+1110.3}_{-981.3}$  arcsec. The innermost part of the profile is nicely fit with by a power law with a slope  $\alpha \sim 0.5$  as typically expected for post-core collapse clusters. Because of the double fit performed in this work, results cannot be directly compared with the literature, where fits have only been performed by using a single King model.



**Fig. 9.** Surface brightness profile for NGC 6254. *Euclid* data are shown as black, filled symbols, while literature values are represented by open symbols. The black, solid line marks the best-fit King model, and the grey shaded area indicates the associated uncertainty. Finally, the bottom panel shows the residuals of the fit.

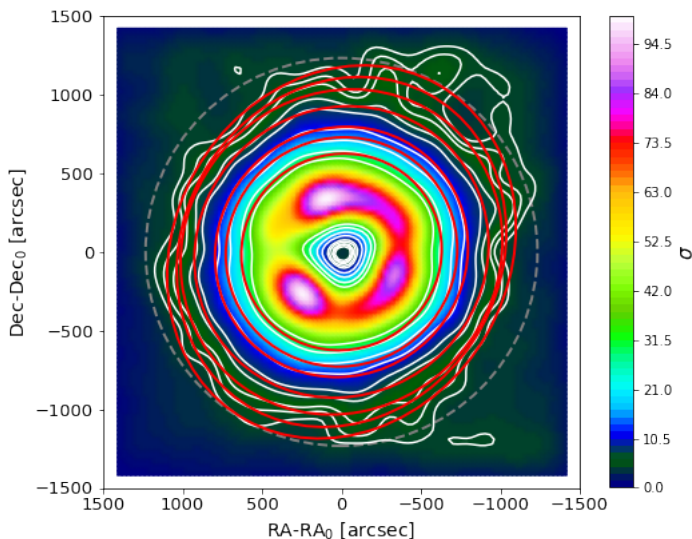


**Fig. 10.** Same as in Fig. 9 but for NGC 6397.



#### 4.2. Outer regions morphology

Given the outstanding performance shown by our ERO data in terms of surface-brightness depth, we use the same sample of member stars described above to look for possible evidence of tidal tails in the two-dimensional morphology of NGC 6254 and NGC 6397. To do so we start from NGC 6254 and compute the stellar-density map shown in Fig. 11 with over-plotted density contours at regular density intervals.



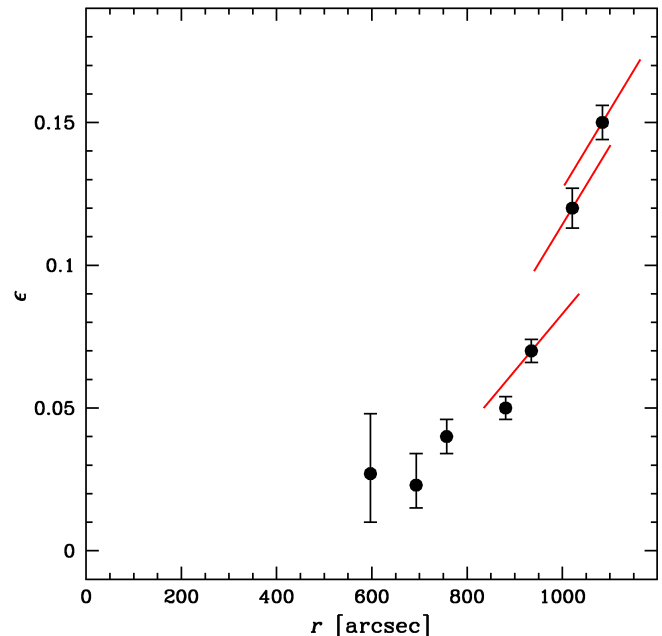
**Fig. 11.** Two-dimensional morphology of NGC 6254 likely member stars. The map is colour-coded based on the stellar density, in units of  $\sigma$  above the background. Iso-density contours are shown in white. The best-fitting ellipses are over-plotted in red. The tidal radius as determined in this work is marked as a dashed grey circle.

As is evident from the figure, the contours are rather circular in the inner regions but, moving outwards, they clearly develop an elongation towards the north-west direction. To provide a quantitative statistical significance of the density contours plotted, we selected members stars in a region more distant than 20 arcmin from the cluster centre and computed the mean stellar density and the dispersion around the mean ( $\sigma_\rho$ ) in the cells of a  $40 \times 40$  grid. The lowest-density contour shown in Fig. 11 corresponds to a density value of  $5.0 \sigma_\rho$ .

We fit the iso-density contours with ellipses to get a quantitative idea of the morphological elongation in the cluster’s outer regions. We fixed the centres of the ellipses on (0,0), and left the angle and the axis ratio as free parameters of the fit. The solutions of the fit are shown in Fig. 11 as red ellipses. The associated uncertainties have been estimated by means of a bootstrapping algorithm, repeated 1000 times. For each iso-density contour, we discarded 50% of the sample of points used to estimate the ellipticity,<sup>4</sup> and performed the fit again. The 16th- and 84th percentiles of the resulting distributions of  $\epsilon$  are taken as asymmetric error estimates. The ellipticity of the iso-density contours is about zero in the inner regions and then progressively increases towards the GC peripheries, as expected in case of tidally induced distortions, reaching a maximum of  $\epsilon = 0.15$ . This trend

<sup>4</sup> defined as  $\epsilon = 1 - (b/a)$ , where  $a$  is the semi-major axis and  $b$  is the semi-minor axis of the ellipse.

is further described in Fig. 12, where the major-axis angles<sup>5</sup>  $\Theta$  of the ellipses for the contours deviating most from circularity are represented by the tilt shown by the solid red lines, namely  $\Theta = 60.5$  deg,  $\Theta = 61.4$  deg and  $\Theta = 51.6$  deg, for a mean angle of  $\langle \Theta \rangle = 57.8$  deg.

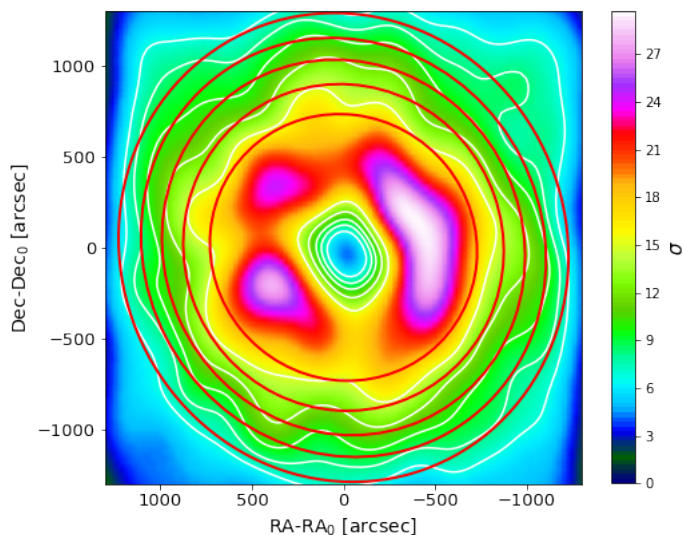


**Fig. 12.** Ellipticity of the iso-density contours as a function of distance from the centre of NGC 6254. The red lines indicate the position angle  $\Theta$  of the best-fitting ellipse.

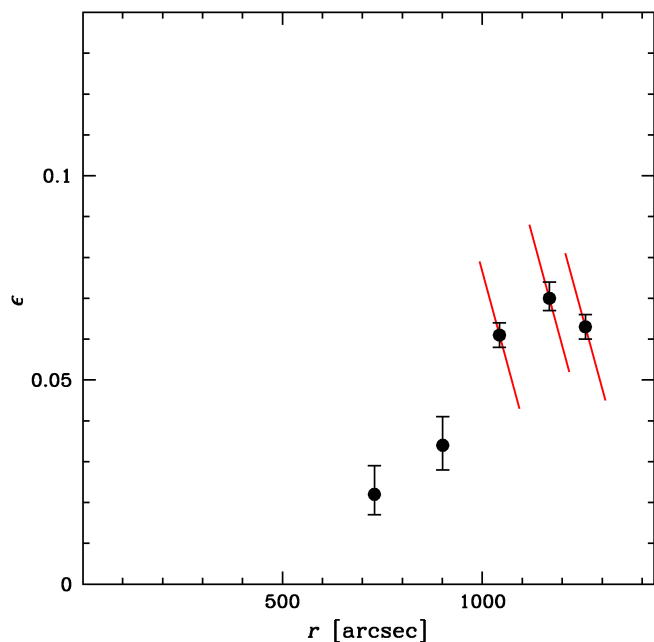
When performing the same analysis on NGC 6397, we obtain the results shown in Fig. 13. The first obvious feature is that the contours appear less elongated compared to the case of NGC 6254. Also in this case they appear circular in the inner regions, and then develop a somewhat higher ellipticity towards the cluster periphery, in the north–south direction.

The solution of the ellipsoidal fitting shown in Fig. 14 confirms the milder morphological elongation. A maximum ellipticity of  $\epsilon = 0.07$  is found in the two most external contours. The fact that we detect a weaker morphological distortion than in NGC 6254 is not surprising. The observed area for NGC 6397 covers a smaller portion of the cluster extent, as NGC 6397 is closer to the Sun and therefore appears more extended on the sky. Our observations can sample only out to about half of the cluster tidal radius of  $r_t \sim 0.8$  deg (see Sect. 4.1). This is not likely far enough to detect distortions related to tidally induced effects. It is nonetheless reassuring that the ellipticity we find matches the one quoted by Harris (1996) and Chen & Chen (2010). The latter work also quotes a direction of the elongation ( $11 \pm 2$  deg oriented anti-clockwise from the north) that is very similar to ours. In fact, we find that the angle of the best-fit ellipses of the three most elongated density contours of our NGC 6397 map is rather stable, varying between  $\Theta = 109.5$  deg and  $\Theta = 114.2$  deg, with a mean value of  $\Theta = 111.0$  deg, corresponding to 21.0 deg in the reference system of Chen & Chen (2010).

<sup>5</sup> defined to increase from west to north, with  $\Theta = 0$  degrees coinciding with the west direction



**Fig. 13.** Two-dimensional morphology of NGC 6397 likely member stars. The map is colour-coded based on the stellar density, in units of  $\sigma$  above the background. Iso-density contours are shown in white. The best-fitting ellipses are over-plotted in red. Note that the tidal radius of NGC 6397 is outside the limits of the figure.



**Fig. 14.** Ellipticity of the iso-density contours as a function of distance from the centre of NGC 6397. The red lines indicate the angle of the best-fitting ellipse.

## 5. Predictions from $N$ -body simulations of NGC 6254

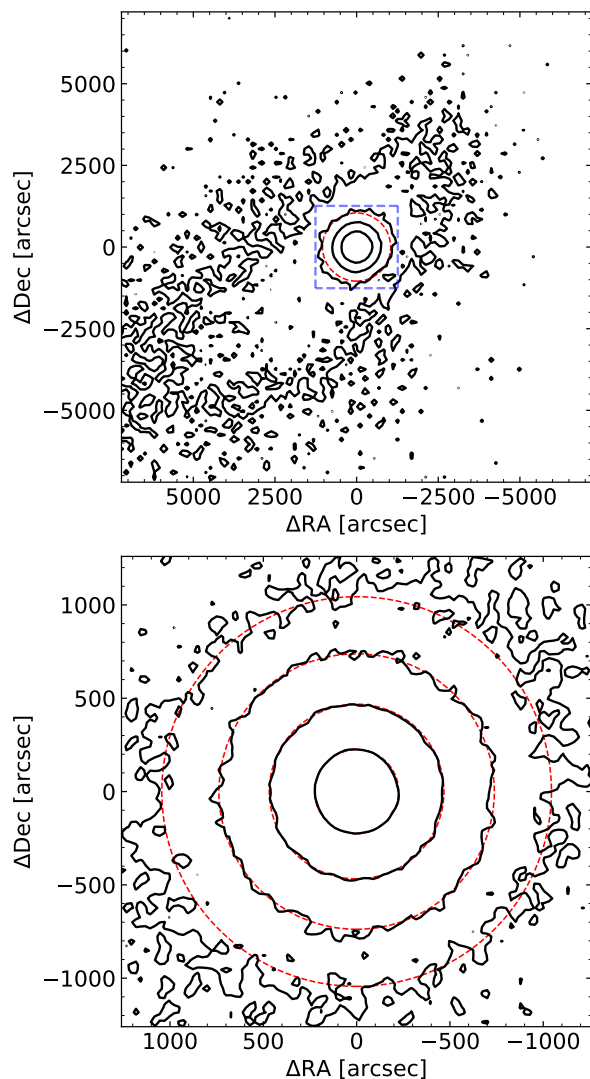
In order to understand how the morphological features observed in NGC 6254 and NGC 6397 compare with expectations, we perform an  $N$ -body simulation of each globular cluster in the Milky Way potential. This is done using `gadget-3`, which is an updated version of `gadget-2` (Springel 2005). For the

Milky Way potential, we take the `MWPotential2014` model from Bovy (2015), which consists of a Navarro–Frenk–White dark matter halo (Navarro et al. 1997), a Miyamoto–Nagai disk (Miyamoto & Nagai 1975), and a broken-power-law bulge. We take the present-day phase-space coordinates for NGC 6254 and NGC 6397 from Vasiliev & Baumgardt (2021). Starting from their present-day coordinates, we rewind a tracer particle for 2 Gyr to generate the initial phase-space coordinates for our models. For NGC 6254, we inject a King profile with a concentration parameter  $W = 7.06$ , a tidal radius of  $r_t = 38.47$  pc, and a mass of  $1.5 \times 10^5 M_\odot$  (de Boer et al. 2019; McLaughlin & van der Marel 2005). We model this with  $5 \times 10^5$  equal mass particles and a particle softening of 0.0063 pc. For NGC 6397, we inject a King profile with  $W = 8.65$ , a tidal radius of  $r_t = 32.16$  pc, and a mass of  $1.13 \times 10^5 M_\odot$  (de Boer et al. 2019; Vitral et al. 2022). We model this with  $10^5$  equal mass particles with a particle softening of 0.0063 pc. Both globular cluster models are then evolved for 2 Gyr to the present day. We stress that this is an initial  $N$ -body model of both systems to study their expected tidal debris (i.e. orientation and ellipticity). As such, these models are initialised with the present-day density profiles for each globular cluster and are disrupted for a short amount of time compared to their age ( $\sim 13$  Gyr). These models are not designed to faithfully capture important effects such as two-body scattering and mass segregation. Accounting for these effects would require starting with a realistic initial mass function, using a collisional  $N$ -body code like `NBODY6` (Aarseth 1999), and evolving the system for the age of the cluster, as was done in e.g. Baumgardt & Hilker (2018).

In Fig. 15 we show mock observations of our model of NGC 6254. We note that the simulated cluster is offset from the correct present-day location by 0.08 kpc in position,  $3.0 \text{ km s}^{-1}$  in velocity and 0.67 degrees in the plane of the sky from the real NGC 6254. As a result, Fig. 15 shows the coordinates relative to the centre of our simulated cluster. The star counts from the  $N$ -body simulation of NGC 6254 display a clear elongation in the north-west direction, which coincides with that found in the observations. When comparing this simulation with the iso-density contours shown in Fig. 11, their similarity is remarkable. This is highlighted even more clearly in Fig. 16, where the predicted (red triangles) and the observed (black circles) variation of NGC 6254 ellipticity as a function of the cluster radius are compared directly. In both cases, the cluster morphology remains about circular in the inner regions and then starts deviating significantly at about 13.5 arcmin. The ellipticity increases rather sharply in both the observations and the simulation by moving farther out and reaches its maximum of  $\epsilon = 0.15$  at  $r \sim 18.5$  arcmin. We note that in the  $N$ -body simulation, the ellipticity is estimated using the eigenvalues of the reduced moment of inertia tensor formed out of the on-sky positions of the particles in circular annuli.

The consistency between the predictions from the  $N$ -body simulations strengthens significantly the interpretation of the morphological distortion we detect in the outer regions of NGC 6254 as an actual feature. This clearly indicates that NGC 6254 has experienced tidal forces due to its interaction with the Milky Way, and that these forces have distorted the cluster morphology, likely leading to the development of tidal tails on larger scales (outside the GC Jacobi radius).

In Fig. 17 we show mock observations of our model of NGC 6397. As with NGC 6254, our  $N$ -body simulation does not exactly match the present-day location and differs by 0.19 kpc,  $12 \text{ km s}^{-1}$ , and 1.1 degrees on the sky with observations of NGC 6397. As a result, the on-sky angles shown in Fig. 17

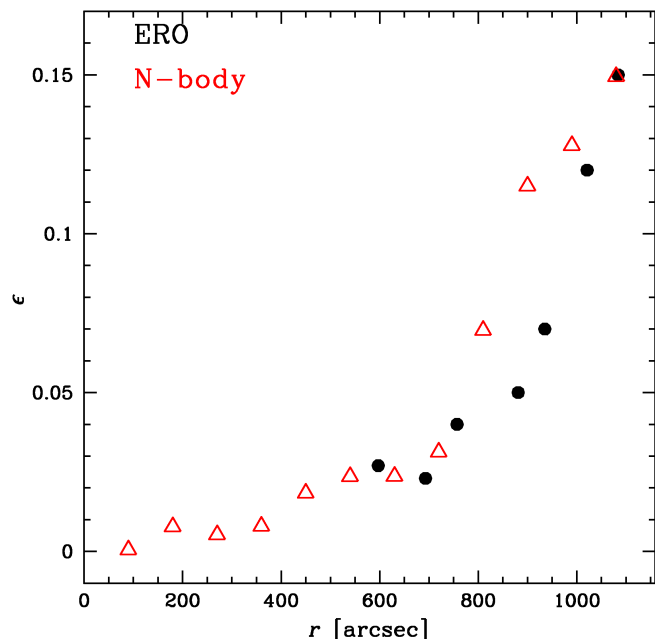


**Fig. 15.** Contour maps of an  $N$ -body simulation of NGC 6254. *Top panel* shows NGC 6254 over a large ( $4 \times 4 \text{ deg}^2$ ) field of view. The dashed-blue square shows the field of view of *Euclid*. A dashed-red circle with a radius of 1044 arcsec is included to show that the simulated cluster is tidally distorted at this radius. The contours are spaced in multiples of 10. *Lower panel* shows the simulation over the field of view of *Euclid* ( $0.7 \times 0.7 \text{ deg}^2$ ). The dashed red circles are chosen to roughly match the contours to show where the cluster is expected to be tidally stretched. The circles are at radii of 225 arcsec, 468 arcsec, 738 arcsec, and 1044 arcsec. The contours are again separated by factors of 10. Note that the largest circle in the bottom panel matches the circle in the top panel.

are relative to the centre of the simulated cluster. Our  $N$ -body simulation of NGC 6397 shows no significant tidal deformation within the *Euclid* field of view, but indicates that tidal deformations should be visible on larger scales. This is consistent with the fact that the Jacobi radius of NGC 6397 is larger than the FoV of *Euclid*.

## 6. Discussion and conclusions

In this paper, we analyse *Euclid* ERO data of two Milky Way GCs, namely NGC 6254 and NGC 6397, to investigate the presence of tidally induced morphological distortions caused by the clusters' interaction with the Milky Way.



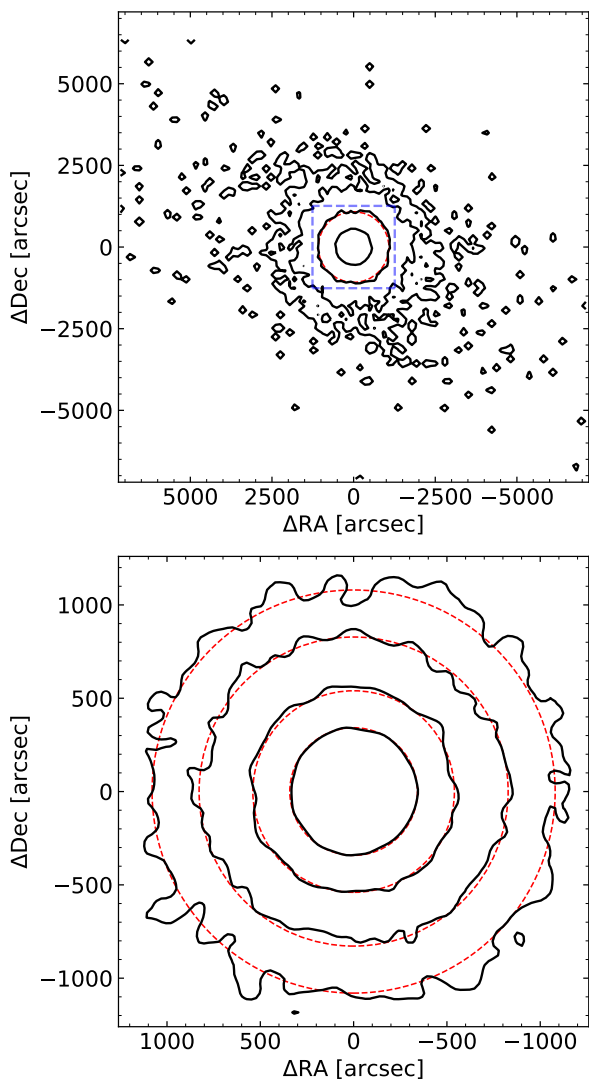
**Fig. 16.** Comparison between the radial variation of the ellipticity observed for NGC 6254 (black symbols) and the prediction from  $N$ -body simulations (red triangles).

As a first result of our analysis, we have shown the first ever *Euclid* CMDs of Galactic GCs. The precision in colour,  $\lesssim 0.04$  mag, seems to indicate the presence of a marginal colour spread in the low-mass regime, that could be associated with GC multiple populations (e.g., Milone et al. 2019) if confirmed by future more detailed studies. The unprecedented combination of wide-field capability and depth further enables the detection of faint, low-mass stars down to  $0.15 M_{\odot}$  and out to the GC tidal radius, at least in the case of NGC 6254. Improved understanding of the data and of the photometric analysis in the future will enable sampling even lower stellar masses.

The surface brightness profiles built from accurately selected samples of cluster members reach more than one magnitude fainter than the surface brightness profiles presented in the literature so far for these two GCs (see e.g., Baumgardt & Hilker 2018), down to a value of  $\mu_V \approx 30.0 \text{ mag arcsec}^{-2}$ . Such a faint limit makes them amongst the deepest ever determined, and the first to be constructed using star counts at the bottom of the MS. The tidal radius estimated from the best-fit King model of NGC 6254 is  $\sim 20\%$  larger than previous estimates, and such an increase can be explained by the fact that *Euclid* enables, for the first time, the use of low-mass MS stars, rather than giants, as morphological tracers.

Most importantly, the findings of this work provide the first robust evidence for tidally induced morphological distortions in the outer regions of NGC 6254, which implies a high chance of finding tidal tails around the cluster on the larger scales, which will be covered by the *Euclid* survey.

This evidence comes from the analysis of the two-dimensional morphology of the cluster. In fact, iso-density contours built from a sample of likely member stars reveal an elongated morphology of NGC 6254 outer regions at more than a  $5\sigma$  significance. Such an elongation is oriented along the north-west direction with a mean angle of about  $58 \text{ deg}$ , which coincides remarkably well with the direction of the cluster tidal



**Fig. 17.** Contour maps of an  $N$ -body simulation of NGC6397. *Top panel* shows NGC 6397 over a large ( $4 \times 4 \text{ deg}^2$ ) field of view. The dashed-blue square shows the field of view of *Euclid*. A dashed-red circle with a radius of 1080 arcsec is included to show that the simulated cluster is not significantly tidally distorted at this radius. The contours are spaced in multiples of 10. *Lower panel* shows the simulation over the field of view of *Euclid*. The dashed-red circles are chosen to roughly match the contours to show where the cluster is expected to be tidally stretched. The circles are at radii of 342 arcsec, 540 arcsec, 828 arcsec, and 1080 arcsec. The contours are separated by half a dex. Note that the largest circle in the bottom panel matches the circle in the top panel.

distortions predicted by  $N$ -body simulations, and disappears towards the cluster inner regions. It should be noted that the direction of the elongation is somewhat reminiscent of the direction in which the extinction varies through the FoV (see Fig. 5). However, our selection of likely cluster members has been performed on a differential reddening-corrected CMD, where the limiting magnitude is the one in the  $H_E$  band. The colour-excess variation  $\delta E(B-V)$  across the field amounts to 0.15 mag corner-to-corner, which based on the extinction law by Cardelli et al. (1989) corresponds to a difference in depth of only 0.08 mag in  $H_E$ . We can thus safely conclude that the detected morphological distortion is not induced by differential extinction-related effects. Moreover, we presented a remarkable similarity between the ellipticity gradient observed from the data and the one predicted

by the  $N$ -body simulations. Both start from an almost circular morphology in the inner  $\sim 10$  arcmin, and then rapidly increase to the maximum value of  $\epsilon = 0.15$  at a distance of  $\sim 18.5$  arcmin. This further supports the notion that the morphological elongations are indeed due to tidal features.

Finally, the two-dimensional morphology of NGC 6397 could only be sampled out to half of the GC tidal radius, and reveals a milder elongation ( $\epsilon_{\text{max}} = 0.07$ ) along about the north-south direction, in agreement with previous studies. Further investigation on larger scales, sampling regions further out than the cluster tidal radius ( $r_t \approx 0.8 \text{ deg}$  as found in Sect. 4.1), is required to more robustly interpret this feature in terms of tidally induced morphological distortion.

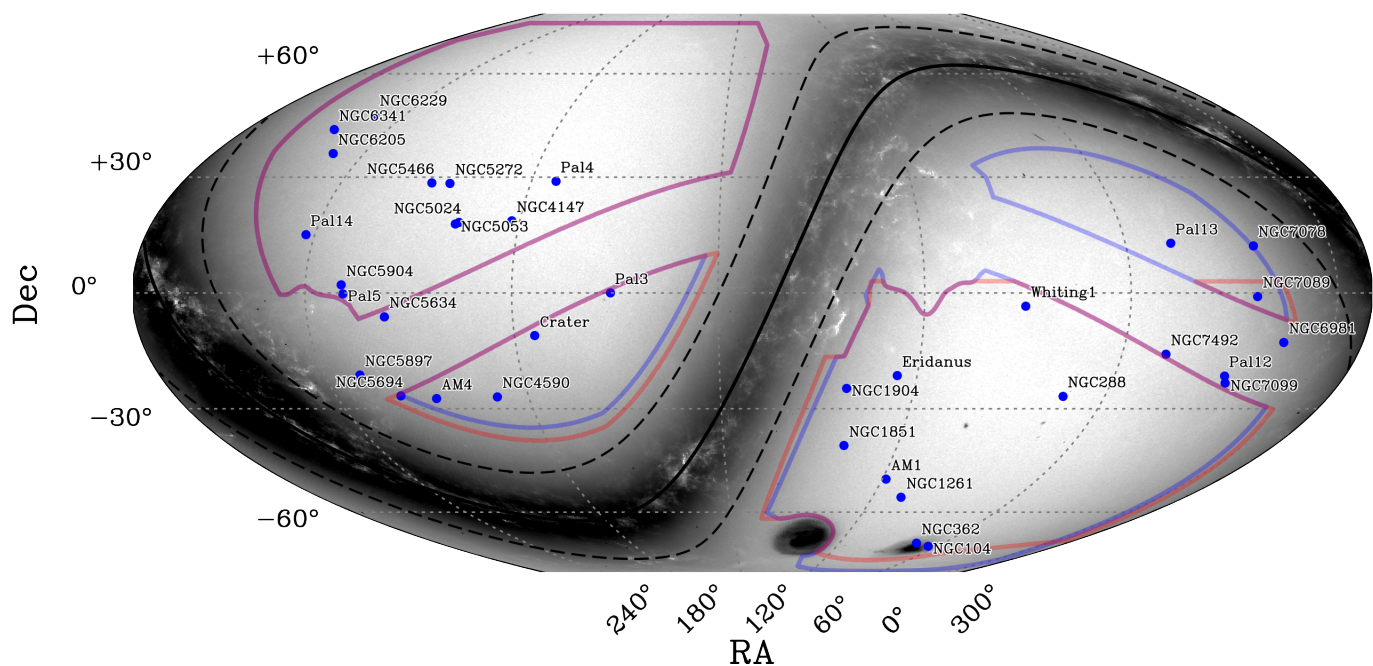
To conclude, the findings of this *Euclid* ERO project demonstrate the potential that the *Euclid* survey has for advancing the study of star clusters in our Galaxy. To determine which GCs fall within the nominal *Euclid* survey coverage, we matched the 6 year *Euclid* Wide Survey Region of Interest (Euclid Collaboration: Scaramella et al. 2022) and planned footprint<sup>6</sup>, shown in Fig. 18 as blue and red polygons, respectively, with a table of the known MW GCs from Baumgardt & Hilker (2018). In total, we find that more than  $\sim 20$  MW GCs will fall in the planned *Euclid* survey footprint. For all of these, the search for tidal tails, as well as the investigation of many other GC-related science cases, can be carried out successfully with *Euclid*.

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<sup>6</sup> We remark that the footprint currently shown in Fig. 18 could be subject to minor changes



**Fig. 18.** Planned Region of Interest (blue polygons) and footprint (red polygons) of the full 6 year *Euclid* wide survey shown in an McBryde-Thomas Flat Polar Quartic projection within blue boxes. The background represents *Gaia* DR3 number counts in logarithmic scale. The position and the names of the GCs covered by the survey or located in its proximity are highlighted in blue.

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